# Complete Solutions Manual to Accompany 

# A First Course in Differential Equations with Modeling Applications <br> ELEVENTH EDITION, METRIC VERSION 

## And

# Differential Equations with Boundary-Value Problems 

# NINTH EDITION, METRIC VERSION 

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## Chapter 1

## Introduction to Differential Equations

### 1.1 Definitions and Terminology

1. Second order; linear
2. Third order; nonlinear because of $(d y / d x)^{4}$
3. Fourth order; linear
4. Second order; nonlinear because of $\cos (r+u)$
5. Second order; nonlinear because of $(d y / d x)^{2}$ or $\sqrt{1+(d y / d x)^{2}}$
6. Second order; nonlinear because of $R^{2}$
7. Third order; linear
8. Second order; nonlinear because of $\dot{x}^{2}$
9. Writing the differential equation in the form $x(d y / d x)+y^{2}=1$, we see that it is nonlinear in $y$ because of $y^{2}$. However, writing it in the form $\left(y^{2}-1\right)(d x / d y)+x=0$, we see that it is linear in $x$.
10. Writing the differential equation in the form $u(d v / d u)+(1+u) v=u e^{u}$ we see that it is linear in $v$. However, writing it in the form $\left(v+u v-u e^{u}\right)(d u / d v)+u=0$, we see that it is nonlinear in $u$.
11. From $y=e^{-x / 2}$ we obtain $y^{\prime}=-\frac{1}{2} e^{-x / 2}$. Then $2 y^{\prime}+y=-e^{-x / 2}+e^{-x / 2}=0$.
12. From $y=\frac{6}{5}-\frac{6}{5} e^{-20 t}$ we obtain $d y / d t=24 e^{-20 t}$, so that

$$
\frac{d y}{d t}+20 y=24 e^{-20 t}+20\left(\frac{6}{5}-\frac{6}{5} e^{-20 t}\right)=24 .
$$

13. From $y=e^{3 x} \cos 2 x$ we obtain $y^{\prime}=3 e^{3 x} \cos 2 x-2 e^{3 x} \sin 2 x$ and $y^{\prime \prime}=5 e^{3 x} \cos 2 x-12 e^{3 x} \sin 2 x$, so that $y^{\prime \prime}-6 y^{\prime}+13 y=0$.
14. Writing $\cos ^{2} x=\frac{1}{2}(1+\cos 2 x)$ and applying $D\left(D^{2}+4\right)$ to the differential equation we obtain

$$
D\left(D^{2}+4\right)\left(D^{2}+4\right)=D\left(D^{2}+4\right)^{2}=0 .
$$

Then

$$
y=\underbrace{c_{1} \cos 2 x+c_{2} \sin 2 x}_{y_{c}}+c_{3} x \cos 2 x+c_{4} x \sin 2 x+c_{5}
$$

and $y_{p}=A x \cos 2 x+B x \sin 2 x+C$. Substituting $y_{p}$ into the differential equation yields

$$
-4 A \sin 2 x+4 B \cos 2 x+4 C=\frac{1}{2}+\frac{1}{2} \cos 2 x .
$$

Equating coefficients gives $A=0, B=1 / 8$, and $C=1 / 8$. The general solution is

$$
y=c_{1} \cos 2 x+c_{2} \sin 2 x+\frac{1}{8} x \sin 2 x+\frac{1}{8} .
$$

59. Applying $D^{3}$ to the differential equation we obtain

$$
D^{3}\left(D^{3}+8 D^{2}\right)=D^{5}(D+8)=0
$$

Then

$$
y=\underbrace{c_{1}+c_{2} x+c_{3} e^{-8 x}}_{y_{c}}+c_{4} x^{2}+c_{5} x^{3}+c_{6} x^{4}
$$

and $y_{p}=A x^{2}+B x^{3}+C x^{4}$. Substituting $y_{p}$ into the differential equation yields

$$
16 A+6 B+(48 B+24 C) x+96 C x^{2}=2+9 x-6 x^{2}
$$

Equating coefficients gives

$$
\begin{aligned}
16 A+6 B & =2 \\
48 B+24 C & =9 \\
96 C & =-6 .
\end{aligned}
$$

Then $A=11 / 256, B=7 / 32$, and $C=-1 / 16$, and the general solution is

$$
y=c_{1}+c_{2} x+c_{3} e^{-8 x}+\frac{11}{256} x^{2}+\frac{7}{32} x^{3}-\frac{1}{16} x^{4} .
$$

60. Applying $D(D-1)^{2}(D+1)$ to the differential equation we obtain

$$
D(D-1)^{2}(D+1)\left(D^{3}-D^{2}+D-1\right)=D(D-1)^{3}(D+1)\left(D^{2}+1\right)=0 .
$$

Then

$$
y=\underbrace{c_{1} e^{x}+c_{2} \cos x+c_{3} \sin x}_{y_{c}}+c_{4}+c_{5} e^{-x}+c_{6} x e^{x}+c_{7} x^{2} e^{x}
$$

57. Solving $\frac{1}{2} q^{\prime \prime}+10 q^{\prime}+100 q=150$ we obtain $q(t)=e^{-10 t}\left(c_{1} \cos 10 t+c_{2} \sin 10 t\right)+3 / 2$. The initial conditions $q(0)=1$ and $q^{\prime}(0)=0$ imply $c_{1}=c_{2}=-1 / 2$. Thus

$$
q(t)=-\frac{1}{2} e^{-10 t}(\cos 10 t+\sin 10 t)+\frac{3}{2} .
$$

As $t \rightarrow \infty, q(t) \rightarrow 3 / 2$.
58. In Problem 54 it is shown that the amplitude of the steady-state current is $E_{0} / Z$, where $Z=\sqrt{X^{2}+R^{2}}$ and $X=L \gamma-1 / C \gamma$. Since $E_{0}$ is constant the amplitude will be a maximum when $Z$ is a minimum. Since $R$ is constant, $Z$ will be a minimum when $X=0$. Solving $L \gamma-1 / C \gamma=0$ for $\gamma$ we obtain $\gamma=1 / \sqrt{L C}$. The maximum amplitude will be $E_{0} / R$.
59. By Problem 54 the amplitude of the steady-state current is $E_{0} / Z$, where $Z=\sqrt{X^{2}+R^{2}}$ and $X=L \gamma-1 / C \gamma$. Since $E_{0}$ is constant the amplitude will be a maximum when $Z$ is a minimum. Since $R$ is constant, $Z$ will be a minimum when $X=0$. Solving $L \gamma-1 / C \gamma=0$ for $C$ we obtain $C=1 / L \gamma^{2}$.
60. Solving $0.1 q^{\prime \prime}+10 q=100 \sin \gamma t$ we obtain

$$
q(t)=c_{1} \cos 10 t+c_{2} \sin 10 t+q_{p}(t)
$$

where $q_{p}(t)=A \sin \gamma t+B \cos \gamma t$. Substituting $q_{p}(t)$ into the differential equation we find

$$
\left(100-\gamma^{2}\right) A \sin \gamma t+\left(100-\gamma^{2}\right) B \cos \gamma t=100 \sin \gamma t
$$

Equating coefficients we obtain $A=100 /\left(100-\gamma^{2}\right)$ and $B=0$. Thus, $q_{p}(t)=\frac{100}{100-\gamma^{2}} \sin \gamma t$. The initial conditions $q(0)=q^{\prime}(0)=0$ imply $c_{1}=0$ and $c_{2}=-10 \gamma /\left(100-\gamma^{2}\right)$. The charge is

$$
q(t)=\frac{10}{100-\gamma^{2}}(10 \sin \gamma t-\gamma \sin 10 t)
$$

and the current is

$$
i(t)=\frac{100 \gamma}{100-\gamma^{2}}(\cos \gamma t-\cos 10 t)
$$

61. In an $L C$-series circuit there is no resistor, so the differential equation is

$$
L \frac{d^{2} q}{d t^{2}}+\frac{1}{C} q=E(t)
$$

Then $q(t)=c_{1} \cos (t / \sqrt{L C})+c_{2} \sin (t / \sqrt{L C})+q_{p}(t)$ where $q_{p}(t)=A \sin \gamma t+B \cos \gamma t$. Substituting $q_{p}(t)$ into the differential equation we find

$$
\left(\frac{1}{C}-L \gamma^{2}\right) A \sin \gamma t+\left(\frac{1}{C}-L \gamma^{2}\right) B \cos \gamma t=E_{0} \cos \gamma t .
$$

Equating coefficients we obtain $A=0$ and $B=E_{0} C /\left(1-L C \gamma^{2}\right)$. Thus, the charge is

$$
q(t)=c_{1} \cos \frac{1}{\sqrt{L C}} t+c_{2} \sin \frac{1}{\sqrt{L C}} t+\frac{E_{0} C}{1-L C \gamma^{2}} \cos \gamma t
$$

From the last result and using $\nu=3 / 2$ we obtain

$$
\begin{aligned}
3 J_{3 / 2}(x) & =x J_{5 / 2}(x)+x J_{1 / 2}(x) \\
J_{5 / 2}(x) & =\frac{3}{x} J_{3 / 2}(x)-J_{1 / 2}(x) \\
& =\frac{3}{x} \sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x}-\cos x\right)-\sqrt{\frac{2}{\pi x}} \sin x \\
& =\sqrt{\frac{2}{\pi x}}\left[\left(\frac{3}{x^{2}}-1\right) \sin x-\frac{3 \cos x}{x}\right]
\end{aligned}
$$

From the last result and using $\nu=5 / 2$ we obtain

$$
\begin{aligned}
5 J_{5 / 2}(x) & =x J_{7 / 2}(x)+x J_{3 / 2}(x) \\
J_{7 / 2}(x) & =\frac{5}{x} J_{5 / 2}(x)-J_{3 / 2}(x) \\
& =\frac{5}{x} \sqrt{\frac{2}{\pi x}}\left(\frac{3 \sin x}{x^{2}}-\frac{3 \cos x}{x}-\sin x\right)-\sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x}-\cos x\right) \\
& =\sqrt{\frac{2}{\pi x}}\left[\left(\frac{15}{x^{3}}-\frac{6}{x}\right) \sin x-\left(\frac{15}{x^{2}}-1\right) \cos x\right]
\end{aligned}
$$

33. (a) To find the spherical Bessel functions $j_{1}(x)$ and $j_{2}(x)$ we use the first formula in (30),

$$
j_{n}(x)=\sqrt{\frac{\pi}{2 x}} J_{n+1 / 2}
$$

with $n=1$ and $n=2$,

$$
j_{1}(x)=\sqrt{\frac{\pi}{2 x} J_{3 / 2}(x)} \quad \text { and } \quad j_{2}(x)=\sqrt{\frac{\pi}{2 x}} J_{5 / 2}(x)
$$

Then from Problem 32 we have

$$
J_{3 / 2}(x)=\sqrt{2} \pi x\left(\frac{\sin x}{x}-\cos x\right) \quad \text { so } \quad j_{1}(x)=\frac{\sin x}{x^{2}}-\frac{\cos x}{x}
$$

and

$$
J_{5 / 2}(x)=\sqrt{2} \pi x\left(\frac{3 \sin x}{x^{2}}-\frac{3 \cos x}{x}-\sin x\right) \quad \text { so } \quad j_{2}(x)=\left(\frac{3}{x^{3}}-\frac{1}{x}\right) \sin x-\frac{3 \cos x}{x^{2}}
$$

(b) Using a graphing utility to plot the graphs of $j_{1}(x)$ and $j_{2}(x)$, we get the red and blue graphcs in the figure to the right.

31. $f(t)=2-2 \mathscr{U}(t-2)+[(t-2)+2] \mathscr{U}(t-2)=2+(t-2) \mathscr{U}(t-2)$
$\mathscr{L}\{f(t)\}=\frac{2}{s}+\frac{1}{s^{2}} e^{-2 s}$
$\mathscr{L}\left\{e^{t} f(t)\right\}=\frac{2}{s-1}+\frac{1}{(s-1)^{2}} e^{-2(s-1)}$
32. $f(t)=t-t \mathscr{U}(t-1)+(2-t) \mathscr{U}(t-1)-(2-t) \mathscr{U}(t-2)=t-2(t-1) \mathscr{U}(t-1)+(t-2) \mathscr{U}(t-2)$
$\mathscr{L}\{f(t)\}=\frac{1}{s^{2}}-\frac{2}{s^{2}} e^{-s}+\frac{1}{s^{2}} e^{-2 s}$
$\mathscr{L}\left\{e^{t} f(t)\right\}=\frac{1}{(s-1)^{2}}-\frac{2}{(s-1)^{2}} e^{-(s-1)}+\frac{1}{(s-1)^{2}} e^{-2(s-1)}$
33. The graph of

$$
f(t)=-1+2 \sum_{k=1}^{\infty}(-1)^{k+1} \mathscr{U}(t-k)=-1+2 \mathscr{U}(t-1)-2 \mathscr{U}(t-2)+2 \mathscr{U}(t-3)-\cdots
$$

is


One way of proceeding to find the Laplace transform is to take the transform term-by-term of the series:

$$
\mathscr{L}\{f(t)\}=-\frac{1}{s}+\frac{2}{s} e^{-s}-\frac{2}{s} e^{-2 s}+\frac{2}{s} e^{-3 s}-\cdots \quad \longleftarrow \quad \text { geometric series }
$$

For $s>0$,

$$
\begin{aligned}
\mathscr{L}\{f(t)\} & =-\frac{1}{s}+\frac{2}{s}\left[e^{-s}-e^{-2 s}+e^{-3 s}-\cdots\right]=-\frac{1}{s}+\frac{2}{s} \cdot \frac{e^{-s}}{1+e^{-s}} \\
& =\frac{e^{-s}-1}{s\left(1+e^{-s}\right)}
\end{aligned}
$$

Alternatively, since $f$ is a periodic functions it can also be defined by

$$
f(t)=\left\{\begin{array}{ll}
-1, & 0 \leq t<1 \\
1, & 1 \leq t<2
\end{array} \quad \text { where } f(t+2)=f(t)\right.
$$

I. If $\lambda=0$ then $X^{\prime \prime}=0$ and $X(x)=c_{1} x+c_{2}$. Also $Y^{\prime \prime}-Y=0$ and $Y(y)=c_{3} \cosh y+c_{4} \sinh y$ so

$$
u=X Y=\left(c_{1} x+c_{2}\right)\left(c_{3} \cosh y+c_{4} \sinh y\right)
$$

II. If $\lambda=-\alpha^{2}<0$ then $X^{\prime \prime}-\alpha^{2} X=0$ and $Y^{\prime \prime}+\left(\alpha^{2}-1\right) Y=0$. The solution of the first differential equation is $X(x)=c_{5} \cosh \alpha x+c_{6} \sinh \alpha x$. The solution of the second differential equation depends on the nature of $\alpha^{2}-1$. We consider three cases:
(i) If $\alpha^{2}-1=0$, or $\alpha^{2}=1$, then $Y(y)=c_{7} y+c_{8}$ and

$$
u=X Y=\left(c_{5} \cosh \alpha x+c_{6} \sinh \alpha x\right)\left(c_{7} y+c_{8}\right) .
$$

(ii) If $\alpha^{2}-1<0$, or $0<\alpha^{2}<1$, then $Y(y)=c_{9} \cosh \sqrt{1-\alpha^{2}} y+c_{10} \sinh \sqrt{1-\alpha^{2}} y$ and

$$
u=X Y=\left(c_{5} \cosh \alpha x+c_{6} \sinh \alpha x\right)\left(c_{9} \cosh \sqrt{1-\alpha^{2}} y+c_{10} \sinh \sqrt{1-\alpha^{2}} y\right) .
$$

(iii) If $\alpha^{2}-1>0$, or $\alpha^{2}>1$, then $Y(y)=c_{11} \cos \sqrt{\alpha^{2}-1} y+c_{12} \sin \sqrt{\alpha^{2}-1} y$ and $u=X Y=\left(c_{5} \cosh \alpha x+c_{6} \sinh \alpha x\right)\left(c_{11} \cos \sqrt{\alpha^{2}-1} y+c_{12} \sin \sqrt{\alpha^{2}-1} y\right)$.
III. If $\lambda=\alpha^{2}>0$, then $X^{\prime \prime}+\alpha^{2} X=0$ and $X(x)=c_{13} \cos \alpha x+c_{14} \sin \alpha x$. Also,

$$
\begin{aligned}
& Y^{\prime \prime}-\left(1+\alpha^{2}\right) Y=0 \text { and } Y(y)=c_{15} \cosh \sqrt{1+\alpha^{2}} y+c_{16} \sinh \sqrt{1+\alpha^{2}} y \text { so } \\
& \qquad u=X Y=\left(c_{13} \cos \alpha x+c_{14} \sin \alpha x\right)\left(c_{15} \cosh \sqrt{1+\alpha^{2}} y+c_{16} \sinh \sqrt{1+\alpha^{2}} y\right) .
\end{aligned}
$$

16. Substituting $u(x, t)=X(x) T(t)$ into the partial differential equation yields $a^{2} X^{\prime \prime} T-g=X T^{\prime \prime}$, which is not separable.
17. Identifying $A=B=C=1$, we compute $B^{2}-4 A C=-3<0$. The equation is elliptic.
18. Identifying $A=3, B=5$, and $C=1$, we compute $B^{2}-4 A C=13>0$. The equation is hyperbolic.
19. Identifying $A=1, B=6$, and $C=9$, we compute $B^{2}-4 A C=0$. The equation is parabolic.
20. Identifying $A=1, B=-1$, and $C=-3$, we compute $B^{2}-4 A C=13>0$. The equation is hyperbolic.
21. Identifying $A=1, B=-9$, and $C=0$, we compute $B^{2}-4 A C=81>0$. The equation is hyperbolic.
22. Identifying $A=0, B=1$, and $C=0$, we compute $B^{2}-4 A C=1>0$. The equation is hyperbolic.
23. Identifying $A=1, B=2$, and $C=1$, we compute $B^{2}-4 A C=0$. The equation is parabolic.
so

$$
A_{0}=0, \quad A_{1}+B_{1}=0, \quad C_{1}+D_{1}=75
$$

and

$$
A_{n}+B_{n}=0, \quad C_{n}+D_{n}=0, \quad \text { for } \quad n>1
$$

When $r=2$

$$
\begin{aligned}
A_{0}+B_{0} \ln 2 & =\frac{1}{2 \pi} \int_{0}^{2 \pi} 60 \cos \theta d \theta=0 \\
A_{n} 2^{n}+B_{n} 2^{-n} & =\frac{1}{\pi} \int_{0}^{2 \pi} 60 \cos \theta \cos n \theta d \theta= \begin{cases}0, & n>1 \\
60, & n=1\end{cases} \\
C_{n} 2^{n}+D_{n} 2^{-n} & =\frac{1}{\pi} \int_{0}^{\infty} 60 \cos \theta \sin n \theta d \theta=0, \quad n=1,2, \ldots,
\end{aligned}
$$

so

$$
B_{0}=0, \quad 2 A_{1}+\frac{1}{2} B_{1}=60, \quad 2 C_{1}+\frac{1}{2} D_{1}=0,
$$

and

$$
A_{n} 2^{n}+B_{n} 2^{-n}=0, \quad C_{n} 2^{n}+D_{n} 2^{-n}=0, \quad \text { for } \quad n>1
$$

Whe have $A_{0}=0$ and $B_{0}=0$, and solving the nonhomogeneous systems for $n=1$,

$$
\begin{array}{rlrl}
A_{1}+B_{1} & =0 & C_{1}+D_{1} & =75 \\
2 A_{1}+\frac{1}{2} B_{1} & =60 & 2 C_{1}+\frac{1}{2} D_{1} & =0
\end{array}
$$

yields $A_{1}=40, B_{1}=-40, C_{1}=-25$, and $D_{1}=100$. Finally, solving the homogeneous systems

$$
\begin{aligned}
A_{n}+B_{n} & =0 & C_{n}+D_{n} & =0 \\
A_{n} 2^{n}+B_{n} 2^{-n} & =0 & C_{n} 2^{n}+D_{n} 2^{-n} & =0
\end{aligned}
$$

gives $A_{n}=B_{n}=C_{n}=D_{n}=0$ for $n>1$. The solution is then

$$
\begin{aligned}
u(r, \theta) & =\left(A_{1} r+B_{1} r^{-1}\right) \cos \theta+\left(C_{1} r+D_{1} r^{-1}\right) \sin \theta \\
& =\left(4-r-40 r^{-1}\right) \cos \theta+\left(-25 r+100 r^{-1}\right) \sin \theta \\
& =40\left(r-\frac{1}{r}\right) \cos \theta-25\left(r-\frac{4}{r}\right) \sin \theta .
\end{aligned}
$$

14. We solve

$$
\begin{gathered}
\frac{\partial^{2} u}{\partial r^{2}}+\frac{1}{r} \frac{\partial u}{\partial r}+\frac{1}{r^{2}} \frac{\partial^{2} u}{\partial \theta^{2}}=0, \quad 0<\theta<\pi, \quad a<r<b \\
u(a, \theta)=\theta(\pi-\theta), \quad u(b, \theta)=0, \quad 0<\theta<\pi \\
u(r, 0)=0, \quad u(r, \pi)=0, \quad a<r<b
\end{gathered}
$$

By writing the boundary condition $x=0$ as

$$
u(0, t)=u_{0}-u_{0} \mathscr{U}(t-1)
$$

its transform is

$$
\begin{aligned}
U(0, s) & =\frac{u_{0}}{s}-\frac{u_{0}}{s} e^{-s} \\
c_{1} & =\frac{u_{0}}{s}-\frac{u_{0}}{s} e^{-s} \\
U(x, s) & =u_{0} \frac{e^{-\sqrt{s} x}}{s}-u_{0} \frac{e^{-\sqrt{s} x}}{s} e^{-s} \\
u(x, t) & =u_{0} \mathscr{L}^{-1}\left\{\frac{e^{-\sqrt{s} x}}{s}\right\}-u_{0} \mathscr{L}^{-1}\left\{\frac{e^{-\sqrt{s} x}}{s} e^{-s}\right\}
\end{aligned}
$$

by entry 3 of Table 14.1.1 and the inverse form of the second translation theorem that:

$$
u(x, t)=u_{0} \operatorname{erfc}\left(\frac{x}{2 \sqrt{t}}\right)-u_{0} \operatorname{erfc}\left(\frac{x}{2 \sqrt{t-1}}\right) \mathscr{U}(t-1)
$$

or

$$
u(x, t)= \begin{cases}u_{0} \operatorname{erfc}\left(\frac{x}{2 \sqrt{t}}\right), & 0<t<1 \\ u_{0} \operatorname{erfc}\left(\frac{x}{2 \sqrt{t}}\right)-u_{0} \operatorname{erfc}\left(\frac{x}{2 \sqrt{t-1}}\right), & t>1\end{cases}
$$

18. The Laplace transform with respect to $t$ of the partial differential equation gives

$$
\frac{d^{2} U}{d x^{2}}-s U=-50 \quad \text { so } \quad U(x, s)=c_{1} e^{-\sqrt{s} x}+c_{2} e^{\sqrt{s} x}+\frac{50}{s}
$$

The boundary condition

$$
\lim _{x \rightarrow \infty} u(x, t)=50 \quad \text { implies } \quad \lim _{x \rightarrow \infty} U(x, s)=\frac{50}{s}
$$

so we take $c_{2}=0$. Thus

$$
U(x, s)=c_{1} e^{-\sqrt{s} x}+\frac{50}{s}
$$

The transform of the boundary condition at $x=0$ is

$$
U(0, s)=\frac{100}{s} e^{-5 s}-\frac{100}{s} e^{-10 s}
$$

Since

$$
\frac{100}{s} e^{-5 s}-\frac{100}{s} e^{-10 s}=c_{1}+\frac{50}{s}
$$

we have

$$
c_{1}=-\frac{50}{s}+\frac{100}{s} e^{-5 s}-\frac{100}{s} e^{-10 s}
$$

